

CAVITATION IN LUBRICATION: AN EVOLUTION MODEL

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1. Formulation.

In this communication we present some of the results of [2] on a model related to the lubrication with cavitation arising in bearings. To describe this phenomenon two unknowns are required: the pressure, p, and the relative content, v. of the oil film. When the lubrication takes an incompressible fluid with convection effects, the Elrod-Adams model to the following equation valid both in the cavitated and the ([13]) leads noncavitated regions

$$(h\gamma)$$
 - $div(h^3\nabla p + h\gamma V) = 0$ in Q,

where $\Omega = (0,T)x\Omega$. $\Omega \subset \mathbb{R}^2$ is a connected open set with regular boundary $\partial\Omega$, $h(t,x,y) \in C^{\infty}(0)$ is a given function with $0 < m \le h \le M$, and V is the given convection term (Some references on the mathematical treatment of this and others rebated models are [3] [4] [6] [7] [9] [12] and [17]). This can be formulated under weak regularity in the following way:

Weak Formulation. Find $(p,y) \in L^2(0,T;H^1(\Omega)) \times L^{\infty}(0)$ such that:

 $p \ge 0$ and $y \in H(p)$ a.e. in Q (here H is the Heaviside graph).

$$\begin{split} &\mathrm{ii)} \quad \int_Q \mathrm{h} \gamma \xi_t = & \int_Q (\mathrm{h}^3 \nabla p \nabla \xi \ + \ \mathrm{h} \gamma V \nabla \xi) \ - \ \int_\Omega \mathrm{h}(0) \gamma_0 \xi(0), \\ & \forall \xi \in & \mathrm{H}^1(Q) \ \ \textit{with} \ \ \xi = 0 \ \ \textit{on} \ \ \partial_1 Q = ((0,T) x \partial \Omega) \cup (\{T\} x \overline{\Omega}). \end{split}$$

iii) p=p on $(0,T)x\partial\Omega$.

2. Existence of Solution by Elliptic Regularization.

We consider the approximated problem:

Elliptic Regularized Formulation. Find p^e \in H¹(Q) such that:

$$i^{\mathcal{E}}$$
) $p^{\mathcal{E}} \ge 0$ a.e.in Q.

$$ii^{\ell}) \quad \epsilon(h^{3}p_{\downarrow}^{\ell})_{\downarrow} + (hF_{\epsilon}(p^{\ell}))_{\downarrow} - div(h^{3}\nabla p^{\ell}) - div(hF_{\epsilon}(p^{\ell})V) = 0 \quad in \ Q$$

iii $^{\mathcal{E}}$) suitable boundary conditions, where \mathbf{F}_{g} is a smooth function approaching the Heaviside graph.

p^e by means of a suitable In [2], we obtain a priori estimates for election of F^E. Then, we show the convergence:

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Finally, p and y satisfy the integral equation ii) and condition i) holds in the sense $\int_{Q} hp(1-y) = 0$. (Similar results in the literature can be found in [15],[1],[11]).

3. Uniqueness.

We shall use, now, the following approach:

Parabolic Regularized Formulation. Find $p^{\varepsilon} \in L^{2}(0,T;H^{1}(\Omega))$ such that:

$$a^{\varepsilon}$$
) $p^{\varepsilon} \ge 0$ a.e.in O.

$$b^{\ell}$$
) $(hF_{\rho}(p^{\ell})) - div(h^{3}\nabla p^{\ell}) - div(hF_{\rho}(p^{\ell})V) = 0$ in Q

$$c^{\ell}$$
) $p^{\ell}=z_{p}$ on $(0,T)\times\partial\Omega$ and $F_{p}(p^{\ell}(0,.))=\gamma_{0}^{\ell}(.)$ on Ω ,

where $p_0^{\varepsilon} > \varepsilon$, $z_{\varepsilon} > \varepsilon$ and F_{ε} is Lipschitz continuous, such that $F_{\varepsilon} \longrightarrow H$ in the sense of graphs, and $F_{\varepsilon}(C(\varepsilon)/2) \ge 1$, (here $p_{\varepsilon} \ge C(\varepsilon)$ with $C(\varepsilon)$ independent on the explicit definition of F_{ε}).

It is not difficult to show that

$$p^{\mathcal{E}} \longrightarrow \overline{p}$$
 (weakly) in $L^2(0,T;H^1(\Omega))$

$$F_o(p^{\mathcal{E}}) \longrightarrow \overline{\gamma}$$
 (weakly star) in $L^{\infty}(Q)$.

with $(\overline{p}, \overline{p})$ solution of the weak formulation which we shall denote in the following as limit solution.

Theorem 1. Let (\hat{p},\hat{y}) be any solution of the weak formulation relative to the initial datum $\hat{\gamma}_0$, let (\bar{p},\bar{y}) be the limit solution relative, to $\hat{\gamma}_0$. Then, for all $t \ge 0$ we have

$$\int_{\Omega} h(t) \left[\hat{\gamma}(t) - \hat{\gamma}(t) \right]^{+} dx \leq \int_{\Omega} h(0) \left[\hat{\gamma}_{0} - \hat{\gamma}_{0} \right]^{+} dx$$

In particular, if $\hat{\gamma}_0 \leq \bar{\gamma}_0$ then $\hat{\gamma}(t) \leq \bar{\gamma}(t)$ and so the solution of the weak formulation is unique.

Idea of the Proof: Let $\xi \in C^{\infty}(Q)$, $\xi(t, \cdot) = 0$ in $\partial \Omega$ By integrating in $Q = (0, t) \times \Omega$ for p and p and substracting we found

$$\begin{split} \int_{\Omega} h(t) \xi(t) \left[\hat{\gamma}(t) - F_{\varepsilon}(p^{\varepsilon}) \right] &= \int_{\Omega} h(0) \xi(0) \left[\hat{\gamma}_{0} - \gamma_{0}^{\varepsilon} \right]^{+} + \\ + \int_{0}^{t} \int_{\Omega} \left[\hat{\gamma} - F_{\varepsilon}(p^{\varepsilon}) \right] \left\{ h \xi_{t} + \frac{\hat{p} - p_{\varepsilon}}{\hat{\gamma} - F_{\varepsilon}(p^{\varepsilon})} \operatorname{divh}^{3} \nabla \xi - h V \nabla \xi \right\} - \\ - \int_{0}^{t} \int_{\partial \Omega} h^{3} \frac{\partial \xi}{\partial v} (\hat{p} - p_{\varepsilon}). \end{split}$$

Denote by $A_{\varepsilon}(t,x) = (\hat{p} - p_{\varepsilon})(\hat{p} - F_{\varepsilon}(p^{\varepsilon}))^{-1}$ and let ξ be the unique solution of the adjoint retrograde non-degenerate elliptic problem:

$$h\xi_1 + A_o(t,x)div(h^3\nabla\xi) - h\nabla\nabla\xi = 0$$
 in $Q_i = (0,t)x\Omega$

$$\begin{split} \xi(t) = & \ \, \mathrm{sign}^{\, \, +} \left[\hat{\gamma}(t) - F_{\mathcal{E}}(p^{\mathcal{E}}(t)) \right] \ \, \mathrm{on} \ \, \Omega \quad \, (\mathrm{final \ condition}) \\ \xi = 0 \quad \, \mathrm{on} \quad \, (0,t) x \partial \Omega. \end{split}$$

From the maximum principle we get $0 \le \xi \le 1$ and using suitable barrier functions we prove:

$$\left| \int_{0}^{t} \int_{\partial \Omega} h^{3} \frac{\partial \xi}{\partial v} (\hat{p} - p_{g}) \right| = O(\varepsilon).$$

Finally

$$\int_{\Omega} h(t) \left[\hat{\gamma}(t) - \bar{\gamma}(t) \right]^{+} dx = \lim_{\varepsilon \to 0} \int_{\Omega} h(t) \left[\hat{\gamma}(t) - F_{\varepsilon}(p^{\varepsilon}(t)) \right]^{+} dx = \lim_{\varepsilon \to 0} \int_{\Omega} h(0) \xi(0) \left[\hat{\gamma}_{0} - \gamma_{0}^{\varepsilon} \right] \leq \int_{\Omega} h(0) \left[\hat{\gamma}_{0} - \bar{\gamma}_{0} \right]^{+}.$$

4. Time Semidiscretization.(Implicit Scheme).

We denote by α the maximal monotone graph H^{-1} . By introducing the new unknown $w=h\gamma$, one has

$$w/h \in H(p) \Leftrightarrow p \in \alpha(w/h)$$
.

Then, w satisfies:

where $w_{0}(.) = h(0,.)\gamma_{0}(.)$.

Using an homogeneous discretization in time $(t_n - t_{n-1} = \lambda)$, the above equation becomes

$$\frac{w(t_n)-w(t_{n-1})}{\lambda} - div \left[h^3(t_n)\nabla\alpha\left(\frac{w(t_n)}{h(t_n)}\right) + w(t_n)\nabla(t_n)\right] \ni 0$$

From this implicit scheme we lead to the following family of problems

where f is a given bounded function.

Now we define the family of abstract operators

$$A_{n}(w) = -\operatorname{div}\left[h^{3}_{n}\nabla\alpha\left(\frac{w}{h_{n}}\right) + wV\right]$$

and

$$D(A_n) = \left\{ w \in L^2 : 0 \le \frac{w}{h_n} \le 1, \alpha \left(\frac{w}{h_n} \right) \in H^1(\Omega), Aw \in L^2, \alpha \left(\frac{w}{h_n} \right) = p_e \text{ on } \partial \Omega \right\}.$$

We have

Theorem 2. There exist a positive constant k depending on V such that if $\lambda \leq k$ and $f \in L^{\infty}$ then there exist an unique $w \in D(A_n)$ satisfying $w + \lambda A_n w \ni f$.

Moreover, if w, and w, are solutions for f, and f, respectively then

$$\left\| \left(\mathbf{w}_{\mathbf{i}} - \mathbf{w}_{\mathbf{2}} \right)^{+} \right\|_{\mathbf{L}^{1}(\Omega)} \leq \left\| \left(\mathbf{f}_{\mathbf{i}} - \mathbf{f}_{\mathbf{2}} \right)^{+} \right\|_{\mathbf{L}^{1}(\Omega)}.$$

Idea of the Proof. We approach α for α_g Lipschitz continuous and such that $\alpha_g \to \alpha$ in the sense of graphs. We get existence for the problem associated to α_g by using pseudo-monotone operators ([16]). The above L¹ estimate is now proved, in that case, by multipliying by sign $(w_1 - w_n)$.

The uniqueness for the regularized problems is obtained again by a duality method, generalizing a preliminary result of [18].

Results on the convergence, when λ goes to zero, are given in [2] in two following cases:

- 1) h and V are time-independent. We get strong convergence in $L^1(\Omega)$ by using an abstract result of [8].
- 2) h and V depend on t in a "regular" way. The L¹ convergence is obtained by showing a suitable condition on the resolvent operator [14]).

5. Monotonicity in time of y

Under suitable conditions we prove the monotonicity of the free boundary defined as $\partial[p=0]$.

Theorem 3. The following inequality holds:

$$\operatorname{div}(\operatorname{Vh})(\chi_0 - \gamma) - \operatorname{Vh} \nabla \gamma - h_1(\chi_0 - \gamma) + h\gamma_1 \ge 0$$
 in Q

where $\chi_0 = \chi$ [p>0]. So, if V=0 then $h_1(\chi_0-\gamma) \le h\gamma_1$ and $\gamma_1 \ge 0$ when $h_1 \le 0$.

Idea of the Proof. We take the test function $H_{\varepsilon}(p)\xi$ with $H_{\varepsilon}(p)=\min\{1,p/\varepsilon\}$ and $\xi\in C_{n}^{\infty}$ such that $\sup \xi(.,x)\in (\tau_{n},T-\tau_{n}),\ \tau_{n}>0$. Then,

$$-\int_O h_t^H {}_{\mathcal{E}}(p) \xi \ = \int_O h^3 \left| \nabla p \right|^2 \! H_{\mathcal{E}}^* \xi \ + \int_O h^3 \nabla p \nabla \xi H_{\mathcal{E}}(p) \ - \int_O \operatorname{div}(\mathbb{V}h) H_{\mathcal{E}}(p) \xi$$

and hence

$$-\!\!\int_Q \!\!\!\! h_t \chi_0 \xi \ + \!\!\! \int_Q \!\!\!\! \operatorname{div}(\mathbb{V} h) \chi_0 \xi - \!\!\! \int_Q \!\!\!\! h^3 \nabla p \nabla \xi \ \chi_0 \ \geq 0 \quad \text{for all } \xi.$$

The result follows by substracting

$$-h_t \chi_0 + \operatorname{div}(Vh)\chi_0 + \operatorname{div}(h^3 \nabla p) \ge 0$$

and

$$-(h\gamma)_{t} + \operatorname{div}(\mathbb{V}h\gamma + h^{3}\nabla p) = 0._{m}$$

Remark. We note $H_{g}(p)\xi \notin H^{1}(Q)$ and so the detailed proof avoids the term $[H_{g}(p)]$. The argument is similar to the one used in [5].

REFERENCES

- [1] Alvarez, S.J. "Problemas de Frontera Libre en Teoria de Lubrificacion" Tesis Doctoral, 1986.
- [2] Alvarez, S.J., Carrillo, J. and Diaz, J.I. Article in preparation.

- [3] Bayada, G., Chambat, M. "Sur quelques modifications de la zone de cavitation en Lubrificacion hydrodynamique". Journal de Mecanique thorique et applique, 5 (1986), 703-729.
- [4] Bayada, G., Chambat, M., El Alaoui, M. Variational Formulation of a Finite Algorithms for Cavitation Problems. To appear in Tribology.1990.
- [5] Carrillo, J. "Uniqueness of the solution of the Evolution Dam Problem". In preparation.
- [6] Chipot, M., Luskin, M. "The compressible Reynolds Lubrication equation". In <u>Metastability and Incompletely Posed Problems</u>, S. Antman, J.L. Ericksen, eds. Springer, 1987, 61-75.
- [7] Cimati, G. "Remark on a Free Boundary Problem of the Theory of Lubrication", B.U.M.I. 6 1982, 249-251.
- [8] Crandall, M., Evans, L.C. "On the relation of the Operator ∂/∂s+∂/∂t to Evolution governed by Accretive Operators" Israel Journal of Mathematics. 21 1975.
- [9] Crank, J. Free and moving boundary problems. Clarendon Press . 1984.
- [10] Diaz J.I., Kersner, R. "On a Nonlinear Degenerate Parabolic Equation on Infiltration or Evaporation through a Porous Medium. Journal of Diff. Eq. 69 (3) 1987 (368-403).
- [11] El Alaoui, M. "Sur un Problme Frontier Libre en Mecanique des Films Minces." These-Lyon I. 1986.
- [12] Elliot, C.M., Ockendon, J.R. "Weak and variational methods for moving boundary problems". Pitman, 1982.
- [13] Elrod, G., Adams, M.L. "A Computer Program for Cavitation and Starva tion Problem" In <u>Cavitation and Related Phenomena in Lubrication</u>" D.Dowson and al.,eds Mech. Eng. Publ. Ltd 1975.
- [14] Evans, L.C., Massey, F.J. "A Remark on the Construction of Nonlinear Evolution Operators." Houston Journal of Math. 4,(1), 1978.
- [15] Gilardi, G. "A New Approach to Evolution Free Boundary Problems". Comm.in P.D.E. 4, (10), 1979.
- [16] Lions, J.L. <u>Quelques methodes de resolution des problemes aux limites</u> non lineaires. Dunod. 1969.
- [17] Rodrigues, J.F. "Obstacle Problems in Mathematical Physics", North-Holland, 1987.
- [18] Rulla, J. "Weak Solutions to Stefan Problems with Prescribed Convection" SIAM J. Math. Anal. 18 1987, 1784-1800.