

# ON THE THERMISTOR PROBLEM WITH TEMPERATURE DEPENDENT CONDUCTIVITY

*Jesús Ildefonso DÍAZ*

*Departamento de Matemática Aplicada  
Universidad Complutense de Madrid, Spain*

Joint work with

**M. Chipot** (Universität Zürich) and **R. Kersner** (Hungarian Academy of  
Sciences, Budapest)

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# 1 Introduction

► **Problem** The one dimensional *thermistor* problem:  $\Omega = (-L, +L)$

$$\begin{cases} u_t - (\kappa(u)u_x)_x = \sigma(u)(v_x)^2 & \text{in } \Omega \times (0, T), \\ (\sigma(u)v_x)_x = 0 & \text{in } \Omega \times (0, T), \\ v = v_D, u = u_D \geq 0 & \text{on } \Gamma_D \times (0, T), \\ \sigma(u)\frac{\partial v}{\partial n} = 0, \kappa(u)\frac{\partial u}{\partial n} = 0 & \text{on } \Gamma_N \times (0, T), \\ u(x, 0) = u_0(x) \geq 0 & \text{on } \Omega. \end{cases} \quad (1)$$

►  $n$  is the outpointing normal vector

►  $\Gamma_D \cup \Gamma_N = \partial\Omega$ ,  $\Gamma_D \cap \Gamma_N = \phi$ , the possibility  $\Gamma_D = \phi$  (the empty set), or  $\Gamma_N = \phi$ , being not excluded.

► This problem models the diffusion of heat produced by Joule's effect in a one dimensional conductor (see for instance Kohlrausch 1900, Cimati 1988).

►  $u$  is the temperature,  $\kappa(u)$  the thermal conductivity of the medium,  $v$  is the inside potential and  $\sigma(u)$  the electric conductivity which (as  $\kappa$  as well) is supposed to depend on the temperature.

► Metallic conduction, the Wiedemann-Franz law  $k(u) = k_0 u \sigma(u) \implies$  the temperature equation degenerates where  $u = 0$ .

► Many results in the literature but: lack of existence and uniqueness results if  $k(0) = 0$  [degenerate equation, finite speed of propagations]. Case of  $\sigma(u)$  degenerate ( $\sigma(0) = 0$ ).

► New formulation:

$$\varphi(s) = \int_0^s \kappa(\tau) d\tau$$

$$\begin{cases} u_t - \varphi(u)_{xx} = \sigma(u)(v_x)^2 & \text{in } \Omega \times (0, T), \\ (\sigma(u)v_x)_x = 0 & \text{in } \Omega \times (0, T), \\ v = v_D, \varphi(u) = \varphi(u_D) \geq 0 & \text{on } \Gamma_D \times (0, T), \\ \sigma(u) \frac{\partial v}{\partial n} = 0, \frac{\partial \varphi(u)}{\partial n} = 0 & \text{on } \Gamma_N \times (0, T), \\ u(x, 0) = u_0(x) \geq 0 & \text{on } \Omega. \end{cases} \quad (2)$$

► No classical solution does not exist in general if  $\kappa(0) = 0$ . Quadratic growth of the right hand side.

► **General assumptions**

$$\sigma \text{ is Lipschitz continuous,} \quad (3)$$

$$\begin{cases} \text{there exists a bounded strictly increasing function } \sigma_0(u), \text{ with } \sigma_0(0) \geq 0 \text{ and } \sigma_1 > 0 \\ \text{such that } \sigma_0(u) \leq \sigma(u) < \sigma_1 \quad \forall u \geq 0, \end{cases} \quad (4)$$

$$\varphi \in C^1([0, +\infty)) \cap C^2((0, +\infty)), \quad (5)$$

$$\varphi'(0) \geq 0, \quad \varphi'(r) > 0 \quad \forall r > 0, \quad (6)$$

$$\begin{cases} \text{there exists } V_D \in L^\infty((0, T); H^1(\Omega)) \text{ such that} \\ V_D = v_D \text{ on } \Gamma_D \times (0, T) \text{ and } \frac{\partial V_D}{\partial n} = 0 \text{ on } \Gamma_N \times (0, T), \end{cases} \quad (7)$$

$$\begin{cases} \text{there exists } U_D \text{ such that } \varphi(U_D) \in H^1((0, T); H^1(\Omega)) \\ \varphi(U_D) = \varphi(u_D) \geq 0 \text{ on } \Gamma_D \times (0, T) \text{ and } \frac{\partial \varphi(U_D)}{\partial n} = 0 \text{ on } \Gamma_N \times (0, T), \end{cases} \quad (8)$$

$$u_0 \in L^\infty(\Omega), \quad 0 \leq u_0 \leq M, \quad (9)$$

► Notice that the Wiedemann-Franz law and the assumption (4) imply

$$k_0 \int_0^u \sigma_0(s) ds \leq \varphi(u) < \tilde{C}u^2 \quad \forall u \geq 0, \quad (10)$$

with  $\tilde{C} = \frac{k_0 \sigma_1}{2}$ .

► Our existence result will require the additional condition

$$\begin{cases} \sigma_0(0) > 0 \\ \text{or} \\ \varphi(u)^\alpha \leq \sigma_0(u) \text{ for any } u \in [0, \delta], \text{ for some } \alpha \in (0, 1) \text{ and } \delta > 0. \end{cases} \quad (11)$$

► The great generality allowed on  $\sigma(u)$  requires to spend some words on the way in which the boundary conditions are satisfied. We shall show that  $\sigma(u)v_x \in L^\infty(Q)$  and that  $\varphi(u(\cdot, t))$  is continuous. Then the assumption

$$\sigma(u_D(x, t)) > 0 \text{ on } \Gamma_D \times [0, T] \quad (12)$$

implies that the trace of  $v$  on  $\Gamma_D \times (0, T)$  is well defined.

► It turns out that a function which plays a crucial role in the study of the system is the function

$$J := \sigma(u)v_x,$$

which corresponds to *the current density*. Notice that the second equation of (1) implies that  $J$  is independent of  $x$ , i.e., for a.e.  $t \in (0, T)$

$$\sigma(u(x, t))v_x(x, t) = J(t) \text{ for a.e. } x \in \Omega. \quad (13)$$

► Since the first equation can be, equivalently written as

$$u_t - \varphi(u)_{xx} = Jv_x \text{ in } \Omega \times (0, T),$$

if  $J(t) \equiv 0$  on some subinterval  $(t_1, t_2) \subset (0, T)$  then the equations of system (1) are not coupled on  $\Omega \times (t_1, t_2)$ .

► Notice also that  $J(t) \equiv 0$  on  $(0, T)$  if  $\inf_{x \in \Omega} |v_x(x, t)| = 0$  (case, for instance, of  $\Gamma_N \neq \emptyset$ ) or  $\min_{x \in \bar{\Omega}} \sigma(u(x, t)) = 0$  (case, for instance, of  $\Gamma_D \neq \emptyset$ ,  $\sigma_0(0) = 0$  and  $u_D(t, x) = 0$ ).

► Moreover, if  $\min_{x \in \bar{\Omega}} \sigma(u(x, t)) > 0$  we have

$$v_x(x, t) = \frac{J(t)}{\sigma(u(x, t))} \text{ a.e. } x \in \Omega.$$

Then, a simple integration shows that, for a.e.  $t \in (0, T)$

$$v(L, t) - v(-L, t) = J(t) \int_{\Omega} \frac{dx}{\sigma(u(x, t))}, \quad (14)$$

which will play an important role in our proof of the existence of solutions and also can be understood as a weak sense in which the Dirichlet condition holds (notice that if  $J(t) \equiv 0$  and, both,  $\Gamma_N \neq \emptyset$  and  $\Gamma_D \neq \emptyset$  then, necessarily,  $v(x, t) = v_D(x, t)$  on  $\Gamma_N \times (0, T)$ ).

► Notice also that if  $J(t) = 0$  and  $\sigma(u(x, t)) > 0$  for any  $x \in \Omega$ , we get that  $v_x \equiv 0$ . Finally, if  $\Gamma_N = \emptyset$  as  $\int_{\Omega} \sigma(u(x, t)) dx > 0$  for a.e.  $t \in (0, T)$ , we get from (14) that  $J(t) = 0$  (respectively  $J(t) > 0$  or  $J(t) < 0$ ) if and only if  $v_D(L, t) - v_D(-L, t) = 0$  (respectively  $v_D(L, t) - v_D(-L, t) > 0$  or  $v_D(L, t) - v_D(-L, t) < 0$ ).

- ▶ The uniqueness of a weak solution will be obtained for the cases in which  $\Gamma_D = \emptyset$  or  $u_D(t, x) > 0$  on  $\Gamma_D \times (0, T)$  (notice that the possible vanishing of  $u_0$  maintains the degenerate character to the parabolic equation).
- ▶ The last section of the paper is devoted to the study of a qualitative property which is peculiar to the case  $\varphi'(0) = 0$ . It concerns with the occurrence of a *free boundary* (given by the boundary of the support of  $u$ ).
- ▶ When  $\sigma_0(0) > 0$  the vanishing set of the solution can be reduced to some curves in  $Q$ .
- ▶ Nevertheless, if  $\sigma_0(0) = 0$ , we show the, so called, *finite speed of propagations* property: if  $u_0(x) = 0$  on  $B_{\rho_0}(x_0) := (x_0 - \rho_0, x_0 + \rho_0)$  for some  $x_0 \in \Omega$  and  $\rho_0 > 0$  then there exists  $t^* > 0$  and a function  $\rho(t) : [0, t^*) \mapsto [0, \infty)$ , with  $\rho(0) \leq \rho_0$ , such that  $u(x, t) = 0$  a.e. in  $B_{\rho(t)}(x_0), \forall t \in [0, t^*)$ .
- ▶ This result opens the possibility of further studies on the properties (and regularity) on this free boundary.

## 2 Existence of a weak solution

Assumed (12), by a **weak solution** to problem (2) we mean a couple of functions  $(u, v)$  such that

$$\varphi(u) \in L^2(0, T; H^1(\Omega)), u \geq 0, u \in C([0, T]; L^1(\Omega)) \cap L^\infty(Q), \quad (15)$$

$$v \in L^\infty(Q), \quad (16)$$

$$\sigma(u)v_x \in L^1(0, T; L^1(\Omega)), \sigma(u)|v_x|^2 \in L^1(0, T; L^1(\Omega)), \quad (17)$$

the boundary conditions  $v = v_D$ ,  $\varphi(u) = \varphi(u_D)$  and  $\sigma(u)\frac{\partial v}{\partial n} = 0, \frac{\partial \varphi(u)}{\partial n} = 0$  hold on  $\Gamma_D \times (0, T)$  and  $\Gamma_N \times (0, T)$  respectively,  $u(\cdot, 0) = u_0$  in  $L^1(\Omega)$  and

$$\begin{aligned} \int_{\Omega} u(x, T)\xi(x, T)dx - \int_{\Omega} u_0(x)\xi(x, 0)dx &= \int_0^T \int_{\Omega} u \xi_t dt dx \\ &- \int_0^T \int_{\Omega} \varphi(u)_x \xi_x dt dx - \int_0^T \int_{\Omega} \sigma(u)|v_x|^2 \xi dt dx, \end{aligned}$$



(18)

$$\int_{\Omega} \sigma(u)v_x \zeta_x dx = 0 \quad \text{a.e. } t \in (0, T),$$

(19)

for all  $\xi, \zeta \in C^1(\overline{Q})$  such that  $\xi(x, t), \zeta(x, t) = 0$  on  $\Gamma_D \times (0, T)$ ..

**Theorem 1.** *Under the assumption (12) there exists, at least, one weak solution to the problem (2). Moreover,  $J(t) := \sigma(u(x, t))v_x(x, t)$  is a bounded (constant in  $x$ ) function on  $(0, T)$  and if  $\sigma_0(0) > 0$  then  $v_x \in L^\infty(Q)$ .*

**Proof.** We can always assume that

$$v_D(L, t) \neq v_D(-L, t) \quad \text{a.e. } t \in (0, T). \quad (20)$$

Indeed, otherwise, as pointed out at the Introduction,  $J(t) \equiv 0$  on  $(0, T)$  and the system is reduced to two uncoupled equations for which the existence of solutions is well-known in the literature.

► Notice that the same appears if  $\Gamma_N = \emptyset$ , nevertheless we shall not assume this condition in the rest of the proof in order to recall an approximation argument which will be used in the proof of the uniqueness.

► The process of proof consists in three steps.

► *Step 1: Approximation.* The method consists in approximating the solution  $(u, v)$  by  $(u^\epsilon, v^\epsilon)$  the solution of

$$\begin{cases} u_t^\epsilon - \varphi(u^\epsilon)_{xx} = \sigma(u^\epsilon)(v_x^\epsilon)^2 & \text{in } \Omega \times (0, T), \\ (\sigma(u^\epsilon)v_x^\epsilon)_x = 0 & \text{in } \Omega \times (0, T), \\ v^\epsilon = v_D, \varphi(u^\epsilon) = \varphi(\max(u_D, \epsilon)) & \text{on } \Gamma_D \times (0, T), \\ \frac{\partial v^\epsilon}{\partial n} = 0, \frac{\partial \varphi(u^\epsilon)}{\partial n} = 0 & \text{on } \Gamma_N \times (0, T), \\ u^\epsilon(\cdot, 0) = u_0 + \epsilon & \text{on } \Omega. \end{cases} \quad (21)$$

► As  $u_t^\epsilon - \varphi(u^\epsilon)_{xx} \geq 0$ , we get

$$u^\epsilon \geq \epsilon \text{ a.e. on } \Omega \times (0, T). \quad (22)$$

Thus,  $\varphi'(u^\epsilon) > 0$ , the operator is, now, **uniformly parabolic** and so a solution  $(u^\epsilon, v^\epsilon)$  to problem (21) is known to exist (see, e.g., Cimatti 1988).

*Step 2: A priori estimates.* We show that  $\sigma(u^\epsilon)(v_x^\epsilon)^2$  (respectively  $u^\epsilon$  and  $\sigma(u)v_x \in L^\infty(Q)$ ) is bounded in  $L^1(Q)$  (respectively in  $L^\infty(Q)$ ) independently of  $\epsilon$ . [For that,

multiply the equation of  $v^\epsilon$  in problem (21) by  $v^\epsilon - V_D$ , integrate by parts, applying the Cauchy-Schwarz inequality,

►  $\sigma(u^\epsilon)(v_x^\epsilon)^2$  bounded in  $L^\infty(0, T : L^1(\Omega))$  implies that  $u_t^\epsilon - \varphi(u^\epsilon)_{xx} = f^\epsilon(t, x)$  with  $f^\epsilon$  uniformly bounded in  $L^1(Q)$  and so we have from Kawanago (1993)  $u^\epsilon$  is uniformly bounded in  $L^\infty(Q)$ .

► On the other hand, from the equation of  $v^\epsilon$  we have

$$\sigma(u^\epsilon)v_x^\epsilon = J^\epsilon(t)$$

and hence

$$\frac{|J^\epsilon(t)|}{\sigma_1} \leq |v_x^\epsilon| \quad (23)$$

Plugging this into (??) we obtain

$$|J^\epsilon(t)|^2 \leq C(T)^2 \sigma_1^2 \text{ess sup}_{t \in [0, T]} \left( \frac{1}{\int_\Omega \sigma(u^\epsilon(x, t)) dx} \right) \quad (24)$$

where  $C(T)$  denotes some constant independent of  $\epsilon$  and so

$$|J^\epsilon(t)| \leq C^*(T) \quad (25)$$

for some positive constant independent of  $\epsilon$ .

► It is easy to get a  $L^\infty(Q)$  a priori estimate on  $v^\epsilon$  since, if  $\Gamma_D \neq \phi$  then, by the maximum principle,

$$|v^\epsilon(x, t)| \leq \|v_D\|_{L^\infty(Q)}, \text{ for a.e. } x \in \Omega \text{ and any } t \in [0, T].$$

In the case  $\Gamma_D = \phi$  the function  $v_x^\epsilon(x, t) = 0$  and since  $v^\epsilon$  is determined up a constant we can take  $v^\epsilon(x, t) = 0$  for a.e.  $x \in \Omega$  and any  $t \in [0, T]$ .

► If  $J_0(t) \neq 0$  for any  $t \in [0, T]$

$$\int_0^T \int_\Omega |v_x^\epsilon| dx dt \leq \frac{\int_0^T \int_\Omega \sigma(u^\epsilon) (v_x^\epsilon)^2 dx dt}{\min_{t \in [0, T]} |J^\epsilon(t)|} \leq C(T). \quad (26)$$

To get other *a priori* estimates we see that by multiplying the equation of  $u^\epsilon$  in problem (21) by  $\varphi(u^\epsilon) - \varphi(\max(U_D, \epsilon)) \in L^2(0, T; H^1(\Omega))$  and integrating over  $\Omega$ , we get that if

$$B(s) = \int_0^s \varphi(u) du$$

and integrating over  $(0, t)$  we obtain for some new constants  $\widehat{C}(T)$  and

$$\begin{aligned} & \int_{\Omega} B(u^\epsilon(x, t)) dx + \frac{1}{2} \int_0^t \int_{\Omega} (\varphi(u^\epsilon)_x)^2 dt dx \\ & \leq \widehat{C}(T) + \int_{\Omega} B(u_0) dx, \end{aligned} \quad (27)$$

► Then, since  $B(s) \geq 0$  for  $s \geq 0$ , there exists some constant  $C = C(T)$ , independent of  $\epsilon$ , such that

$$\begin{aligned} & \int_{\Omega} B(u^\epsilon(., t)) \leq C(T) \quad \forall t \in (0, T), \\ & \|\varphi(u^\epsilon)\|_{L^2(0, T; H^1(\Omega))} \leq C(T), \end{aligned}$$

and (from the equation of  $u^\epsilon$ )

$$\|u_t^\epsilon\|_{L^2(0, T; H^{-1}(\Omega))} \leq C(T).$$

Moreover, since

$$\varphi(u^\epsilon)_t = \varphi'(u^\epsilon)u_t^\epsilon,$$

it follows that

$$\|\varphi(u^\epsilon)_t\|_{L^2(0, T; H^{-1}(\Omega))} \leq C(T).$$

*Step 3 : Passage to the limit.* Using a classical compactness argument, (see [30]) and the monotonicity of  $\varphi$ , we can extract a “subsequence” that for simplicity we still label by “ $\epsilon$ ” such that

$$\varphi(u^\epsilon) \rightharpoonup l_1 \text{ in } L^2(0, T; H^1(\Omega)), \quad (28)$$

$$\varphi(u^\epsilon) \rightarrow l_1 \text{ in } L^2(Q), \quad (29)$$

$$u^\epsilon \rightarrow u \text{ in } L^\infty(Q), \quad (30)$$

$$(u^\epsilon)_t \rightharpoonup l_2 \text{ in } L^2(0, T; H^{-1}(\Omega)), \quad (31)$$

$$J^\epsilon(t) \rightharpoonup J(t) \text{ weakly-star in } L^\infty(0, T), \quad (32)$$

$$v^\epsilon \rightharpoonup v \text{ weakly-star in } L^\infty(0, T : L^\infty(\Omega)), \quad (33)$$

$$\sigma(u^\epsilon) (v_x^\epsilon)^2 = J^\epsilon(t)v_x^\epsilon \rightharpoonup l_3 \text{ weakly-star in } L^\infty(0, T : L^1(\Omega)). \quad (34)$$

► Clearly, one deduces that  $l_1 = \varphi(u)$ ,  $l_2 = u_t$ , . In order to prove the regularity  $u \in C([0, T]; L^1(\Omega))$  it suffices to multiply the equation of  $u^\epsilon$  by  $\text{sign}(\varphi(u^\epsilon))$ . Then,

$$\frac{d}{dt} \int_{\Omega} |u^\epsilon(x, t)| dx \leq \int_{\Omega} |\sigma(u^\epsilon(x, t))(v_x^\epsilon(x, t))^2| dx$$

and in the limit

$$\frac{d}{dt} \int_{\Omega} |u(x, t)| dx \leq \int_{\Omega} |l_3(x, t)| dx$$

which proves that  $u \in C([0, T]; L^1(\Omega))$ .

► Since  $\min_{x \in \bar{\Omega}} \sigma(u^\epsilon(x, t)) > 0$ , by using the identity

$$v_D(L, t) - v_D(-L, t) = J^\epsilon(t) \int_{\Omega} \frac{dx}{\sigma(u^\epsilon(x, t))}, \quad (35)$$

we deduce that, for a.e.  $t \in (0, T)$ ,

$$J^\epsilon(t) \rightarrow J(t) \text{ in } \mathbb{R} \quad (36)$$

since for a.e.  $t \in (0, T)$

$$\frac{v_D(L, t) - v_D(-L, t)}{\int_{\Omega} \frac{dx}{\sigma(u^\epsilon(x, t))}} = J^\epsilon(t), \quad (37)$$

$\sigma(u^\epsilon(x, t)) \rightarrow \sigma(u(x, t))$  for any  $x \in \Omega$  (recall that  $\varphi(u^\epsilon) \rightharpoonup \varphi(u)$  in  $L^2(0, T; H^1(\Omega))$ ) implies that  $\varphi(u^\epsilon(\cdot, t)) \rightarrow \varphi(u(\cdot, t))$  in  $C(\bar{\Omega})$  for a.e.  $t \in (0, T)$ ) and notice that if  $\min_{x \in \bar{\Omega}} \sigma(u(x, t)) = 0$  then  $J(t) = 0$  and (36) is reduced to  $J^\epsilon(t) \rightarrow 0$  in  $\mathbb{R}$ .

► Thus, we conclude that for any  $\xi \in C^1(\bar{Q})$  such that  $\xi(x, t) = 0$  on  $\Gamma_D \times (0, T)$

$$\begin{aligned} \int_{\Omega} \sigma(u^\epsilon(x, t)) |v_x^\epsilon(x, t)|^2 \xi(x, t) dx &= J^\epsilon(t) \int_{\Omega} v_x^\epsilon(x, t) \xi(x, t) dx \\ &= -J^\epsilon(t) \int_{\Omega} v^\epsilon(x, t) \xi_x(x, t) dx. \end{aligned}$$

Using (33), (34) and (36) we can pass to the limit to deduce that

$$\int_{\Omega} \sigma(u^\epsilon(x, t)) |v_x^\epsilon(x, t)|^2 \xi(x, t) dx \rightarrow -J(t) \int_{\Omega} v(x, t) \xi_x(x, t) dx = \int_{\Omega} l_3(x, t) \xi(x, t) dx.$$

► Then, we can pass to the limit in the boundary conditions and in the equations to get that

$$\begin{aligned} \int_{\Omega} u(x, T) \xi(x, T) dx - \int_{\Omega} u_0(x) \xi(x, 0) dx - \int_0^T \int_{\Omega} u \xi_t dt dx \\ + \int_0^T \int_{\Omega} \varphi(u)_x \xi_x dt dx = \int_0^T \int_{\Omega} \sigma(u) (v_x)^2 \xi dt dx, \end{aligned} \quad (38)$$

$$\int_{\Omega} \sigma(u) v_x \zeta_x dx = 0 \text{ a.e. } t \in (0, T), \quad (39)$$

to get the existence result. ■

### 3 Uniqueness of solutions

► Our main idea will consist in proving that any possible weak solution must coincide with the solution constructed in the previous section by using a method that, coming from the Holmgren duality method, it was first adapted to degenerate equations by A.S. Kalashnikov (1979) and then refined in Díaz-Kersner (1993)[16].



**Theorem 2.** *Assume  $\Gamma_D = \phi$  or  $u_D(t, x) > 0$  on  $\Gamma_D \times (0, T)$ . Then problem (2) has a unique weak solution  $(u, v)$  such that  $v \in L^1(0, T; W^{1,1}(\Omega))$  ( $v(\cdot, t)$  being univocally determined in  $\Omega$  unless a constant in the case of  $\Gamma_D = \phi$  and arbitrary on the set  $\{(x, t) \in Q, u(t, x) = 0\}$ ).*

► Before giving the proof of this theorem let us introduce some notation. Let  $(u^\epsilon, v^\epsilon)$  be as before. Let  $(w, z)$  be any weak solution to problem (2). By subtracting and using that

$$\int_0^T \int_\Omega (\sigma(w) (z_x)^2) \xi dt dx = \int_0^T \int_\Omega ((\sigma(w) z z_x)_x) \xi dt dx = - \int_0^T \int_\Omega \sigma(w) z_x z \xi_x dt dx,$$

which is justified since  $\sigma(w) z_x \in L^1(Q)$  and  $z \in L^\infty(Q)$ , we have

$$\begin{aligned} & \int_\Omega (w - u^\epsilon)(x, T) \xi(x, T) dx - \int_\Omega (w - u^\epsilon)(x, 0) \xi(x, 0) dx = \int_0^T \int_\Omega (w - u^\epsilon) \xi_t dt dx \\ & \quad + \int_0^T \int_\Omega (\varphi(w) - \varphi(u^\epsilon)) \xi_{xx} dt dx - \int_0^T \int_\Omega (\sigma(w) z z_x - \sigma(u^\epsilon) v^\epsilon v_x^\epsilon) \xi_x dt dx \\ & \quad + \int_0^T \int_{\Gamma_D} (\varphi(\max(U_D(s, t), \epsilon)) - \varphi(u_D(s, t))) \frac{\partial \xi}{\partial n}(s, t) dt ds + \int_0^T \int_\Omega (\sigma(w) z_x - \sigma(u^\epsilon) v_x^\epsilon) \zeta_x dt dx \end{aligned} \tag{40}$$

for all  $\xi, \zeta \in C^1(\overline{Q}) \cap C([0, T] : C^2(\overline{\Omega}))$  such that  $\xi(x, t), \zeta(x, t) = 0$  on  $\Gamma_D \times (0, T)$ .

► Here the term  $\int_0^T \int_{\Gamma_D} (\varphi(\max(U_D(s, t), \epsilon)) - \varphi(u_D(s, t))) \frac{\partial \xi}{\partial n}(s, t) dt ds$  must be understood in the usual onedimensional integration by parts sense. So if, for instance,  $\Gamma_D = \{-L\} \cup \{L\}$  we have that

$$\begin{aligned} & \int_0^T \int_{\Gamma_D} (\varphi(\max(U_D(s, t), \epsilon)) - \varphi(u_D(s, t))) \frac{\partial \xi}{\partial n}(s, t) dt ds \\ &= \int_0^T [(\varphi(\max(U_D(L, t), \epsilon)) - \varphi(u_D(L, t))) \xi_x(L, t) - (\varphi(\max(U_D(-L, t), \epsilon)) - \varphi(u_D(-L, t))) \xi_x(-L, t)] dt \end{aligned}$$

► Let us denote by  $I$  the left hand side of (40). Thus

$$\begin{aligned} I &= \int_0^T \int_{\Omega} (w - u^\epsilon) (\xi_t + \frac{\varphi(w) - \varphi(u^\epsilon)}{w - u^\epsilon} \xi_{xx} - \frac{\sigma(w) - \sigma(u^\epsilon)}{w - u^\epsilon} z z_x \xi_x + \\ &+ \frac{\sigma(w) - \sigma(u^\epsilon)}{w - u^\epsilon} z_x \zeta_x) dt dx + \int_0^T \int_{\Gamma_D} (\varphi(\max(U_D(s, t), \epsilon)) - \varphi(u_D(s, t))) \frac{\partial \xi}{\partial n}(s, t) dt ds \\ &\quad - \int_0^T \int_{\Omega} (z - v^\epsilon) (\sigma(u^\epsilon) z_x \xi_x - (\sigma(u^\epsilon) v^\epsilon \xi_x)_x + (\sigma(u^\epsilon) \zeta_x)_x) dt dx. \quad (41) \end{aligned}$$

► Let us set

$$A_\epsilon = A_\epsilon(x, t) = \frac{\varphi(w) - \varphi(u^\epsilon)}{w - u^\epsilon}, \quad B_\epsilon = B_\epsilon(x, t) = \frac{\sigma(w) - \sigma(u^\epsilon)}{w - u^\epsilon} z z_x \quad (42)$$

$$C_\epsilon = C_\epsilon(x, t) = \frac{\sigma(w) - \sigma(u^\epsilon)}{w - u^\epsilon} z_x, \quad D_\epsilon = D_\epsilon(x, t) = \sigma(u^\epsilon) z_x \quad (43)$$

$$E_\epsilon = E_\epsilon(x, t) = \sigma(u^\epsilon) v^\epsilon, \quad F_\epsilon = F_\epsilon(x, t) = \sigma(u^\epsilon). \quad (44)$$

Thus (41) reads now :

$$I = \int_0^T \int_\Omega (w - u^\epsilon) \{ \xi_t + A_\epsilon \xi_{xx} - B_\epsilon \xi_x + C_\epsilon \zeta_x \} dt dx + \int_0^T \int_{\Gamma_D} (\varphi(\max(U_D(s, t), \epsilon)) - \varphi(u_D(s, t))) \\ - \int_0^T \int_\Omega (z - v^\epsilon) \{ D_\epsilon \xi_x - (E_\epsilon \xi_x)_x + (F_\epsilon \zeta_x)_x \} dt dx.$$

**Lemma 1.** *There exist three positive constants  $m_\epsilon$ ,  $M_\epsilon$  and  $M^*$  ( $M^*$  independent of  $\epsilon$ ) such that  $m_\epsilon \leq A_\epsilon(x, t) \leq M_\epsilon \quad \forall (x, t) \in Q$ ,  $|B_\epsilon(x, t)|, |C_\epsilon(x, t)| \leq M^* \quad \forall (x, t) \in Q$ . ■*

► Assume that we extend  $A_\epsilon$ ,  $B_\epsilon$ ,  $C_\epsilon$  to the whole  $\mathbb{R}^2$  respectively by  $m_\epsilon, 0, 0$  and denote again these extensions by  $A_\epsilon$ ,  $B_\epsilon$ ,  $C_\epsilon$ . Let  $\rho$  be a function of class  $C^\infty$  with

support in the ball  $B(0, 1)$  of center  $\mathbf{0}$  and radius 1 of  $\mathbb{R}^2$  and such that

$$\int_{B(0,1)} \rho dt dx = 1.$$

Set

$$\rho_n(x, t) = n^2 \rho(nx, nt)$$

and

$$A_\epsilon^n = \rho_n * A_\epsilon, \quad B_\epsilon^n = \rho_n * B_\epsilon, \quad C_\epsilon^n = \rho_n * C_\epsilon,$$

where  $*$  denotes the usual convolution of functions. Clearly, these functions are of class  $C^\infty$  in  $\mathbb{R}^2$ . Moreover, one has

$$m_\epsilon \leq A_\epsilon^n(x, t) \leq M_\epsilon \quad \forall (x, t) \in Q, \forall n \quad (45)$$

$$|B_\epsilon^n|, |C_\epsilon^n| \leq M^* \quad \forall (x, t) \in Q, \forall n. \quad (46)$$

Thus, equation (40) reads now

$$I = \int_0^T \int_\Omega (w - u^\epsilon) \{ \xi_t + A_\epsilon^n \xi_{xx} - B_\epsilon^n \xi_x + C_\epsilon^n \zeta_x \} dt dx + \\ \int_0^T \int_{\Gamma_D} (\varphi(\max(U_D(s, t), \epsilon)) - \varphi(u_D(s, t))) \frac{\partial \xi}{\partial n}(s, t) dt ds$$

$$\begin{aligned}
& + \int_0^T \int_{\Omega} (w - u^\epsilon) \xi_{xx} (A_\epsilon - A_\epsilon^n) dt dx - \int_0^T \int_{\Omega} (w - u^\epsilon) \xi_x (B_\epsilon - B_\epsilon^n) dt dx \\
& + \int_0^T \int_{\Omega} (w - u^\epsilon) \zeta_x (C_\epsilon - C_\epsilon^m) dt dx - \int_0^T \int_{\Omega} (z - v^\epsilon) \{ D_\epsilon \xi_x - (E_\epsilon \xi_x)_x + (F_\epsilon \zeta_x)_x \} dt dx.
\end{aligned} \tag{47}$$

A similar argument must be applied if the coefficients  $D_\epsilon$ ,  $E_\epsilon$  and  $F_\epsilon$  are not bounded (we leave the details to the reader).

Now we construct a “dual system” which plays a crucial role in the proof of Theorem 2.

**Lemma 2.** *There exists a unique smooth solution  $(\xi, \zeta) = (\xi_\epsilon^{n,m}, \zeta_\epsilon^{n,m})$  to the system*

$$\begin{cases} \xi_t + A_\epsilon^n \xi_{xx} - B_\epsilon^n \xi_x + C_\epsilon^n \zeta_x = 0 & \text{in } Q, \\ -(F_\epsilon \zeta_x)_x = (E_\epsilon \xi_x)_x - D_\epsilon \xi_x & \text{in } Q, \\ \zeta = 0, \quad \xi = 0 & \text{on } \Gamma_D \times (0, T), \\ \frac{\partial \zeta}{\partial n} = 0, \quad \frac{\partial \xi}{\partial n} = 0 & \text{on } \Gamma_N \times (0, T), \\ \xi(\cdot, T) = w^m & \text{on } (0, L), \end{cases} \tag{48}$$

where  $w^m \in C_0^\infty(\Omega)$  is such that  $|w^m(x)| \leq 1$  for any  $x \in (-L, L)$  and

$$w^m \rightarrow \text{sign}(w(x, T) - u^\epsilon(x, T)) \quad \text{in } L^2(\Omega), \text{ when } m \rightarrow \infty \tag{49}$$

(here  $\text{sign}$  denotes the  $\text{sign}_0$  function, i.e.,  $\text{sign}(x) = x/|x|$  if  $x \neq 0$ , and 0 if  $x = 0$ ).

Next, we show some estimates which we shall need later.

**Lemma 3.** *Let  $(\xi_\epsilon^{n,m}, \zeta_\epsilon^{n,m})$  be the solution to problem (48). Then, there exists a constant  $C_\epsilon$  independent of  $n$  and  $m$  such that*

$$\|\xi_{\epsilon,xx}^{n,m}\|_{2,Q}, \quad \|\xi_{\epsilon,x}^{n,m}\|_{2,Q}, \quad \|\zeta_{\epsilon,x}^{n,m}\|_{2,Q} \leq C_\epsilon \quad (50)$$

( $\|\cdot\|_{2,Q}$  denotes the usual  $L^2$ -norm on  $L^2(Q)$ ).

**Proof of Theorem 2.** In (40), (47) choose  $(\xi, \zeta) = (\xi_\epsilon^{n,m}, \zeta_\epsilon^{n,m})$  solution to (48) where  $w^m$  satisfies (49). Then expressions (40) and (47) leads to

$$\int_{\Omega} (w - u^\epsilon)(x, T) w^m(x) dx - \int_{\Omega} (w - u^\epsilon)(x, 0) \xi(x, 0) dx \quad (51)$$

$$\begin{aligned} &= \int_0^T \int_{\Gamma_D} (\varphi(\max(U_D(s, t), \epsilon)) - \varphi(u_D(s, t))) \frac{\partial \xi}{\partial n}(s, t) dt ds \\ &+ \int_0^T \int_{\Omega} (w - u_\epsilon) \xi_{xx} (A_\epsilon - A_\epsilon^n) dx dt \\ &- \int_0^T \int_{\Omega} (w - u_\epsilon) \xi_x (B_\epsilon - B_\epsilon^n) dx dt + \int_0^T \int_{\Omega} (w - u_\epsilon) \zeta_x (C_\epsilon - C_\epsilon^n) dx dt. \end{aligned} \quad (52)$$

Notice that the assumption  $u_D > 0$  implies that for  $\epsilon > 0$  small enough we get that  $\max(U_D(s, t), \epsilon) = u_D(s, t)$  and so the first term of the right hand side disappears (this is also the case of  $\Gamma_D = \phi$ ). Then, by passing to the limit (first in  $\epsilon \rightarrow 0$ , then in  $n \rightarrow \infty$  and finally in  $m \rightarrow \infty$ ) we get that  $\int_{\Omega} |w(x, T) - u^\epsilon(x, T)| dx = 0$ . Since  $T$  is arbitrary we get that  $w \equiv u^\epsilon$  and then, obviously,  $z \equiv v$  on  $Q$  if  $\Gamma_D \neq \phi$ .  
 ► When  $\Gamma_D = \phi$  we deduce that, for any  $t \in [0, T]$  there exists a constant  $C(t)$  such that  $z(\cdot, t) - v(\cdot, t) = C(t)$  on  $\Omega$ . On the set  $\{(x, t) \in Q, u(t, x) = 0\}$   $z$  and  $v$  may be different without any consequence on the rest of points of  $Q$ . ■

**Remark 2.** The case  $u_D = 0$  is more delicate since the first term of the right hand side of (51) becomes

$$\int_0^T \int_{\Gamma_D} \varphi(\epsilon) \frac{\partial \xi}{\partial n}(s, t) dt ds$$

and the passing to the limit requires sharper estimates obtained under additional assumptions (see, for instance, Díaz-Kersner 1993 for the case of a single scalar equation). Nevertheless, we conjecture that, as in the scalar case the uniqueness of weak solutions holds also for  $u_D = 0$  and general functions  $\varphi$ .

## 4 On the existence of the free boundary

► The assumption  $\varphi'(0) = 0$  and a suitable growth assumption lead to the existence of a free boundary given as the boundary of the support of the solution. It is the *finite speed of propagation property*: if  $u_0(x) = 0$  on  $B_{\rho_0}(x_0) := (x_0 - \rho_0, x_0 + \rho_0)$  for some  $x_0 \in \Omega$  and  $\rho_0 > 0$  then there exists  $t^* > 0$  and a function  $\rho(t) : [0, t^*) \mapsto [0, \infty)$ , with  $\rho(0) \leq \rho_0$ , such that  $u(x, t) = 0$  a.e. in  $B_{\rho(t)}(x_0), \forall t \in [0, t^*)$ .

► When  $\Gamma_N \neq \phi$  we know that the system becomes uncoupled (see the Introduction) and, so, the criterium for the finite speed of propagation is well known (see, e.g., the surveys Kalashnikov (1987) and Antontsev-Díaz-Shmarev 2002).

► Nevertheless, if  $\Gamma_N = \phi$  and  $\sigma_0(0) > 0$  the vanishing set of the solution can be reduced (at most) to some curves in  $Q$  since, if we assume that  $u(\cdot, t)$  is a convex function of  $x$  then

$$u_t \geq \sigma_0(v_x)^2$$

and thus

$$u(x, t) \geq \sigma_0 \int_0^t v_x(x, s)^2 ds + u_0(x).$$

Then  $\int_0^t v_x(x_0, s)^2 ds > 0$  implies that  $u(x_0, t) > 0$ . Notice also that, by the strong maximum principle,  $v_x(x, \cdot)$  can not be zero on a subset of  $\Omega$  of positive measure



(for any fixed  $t \in [0, T]$ ).

► On the other hand, if  $\sigma(u(x, t)) (v_x(x, t))^2 > 0$  on  $Q$ , it is impossible to get solutions  $u(x, t)$  vanishing on an open subset  $\omega$  of  $Q$  since we would reach a contradiction on  $\omega$  through the equation of (1).

► The case  $\sigma_0(0) = 0$  (and  $\Gamma_N = \emptyset$ ) is different. More precisely we have:

**Theorem 3.** *Assume  $\varphi$  satisfying*

$$\int_{0^+} \frac{\varphi'(s)}{s} ds < \infty, \quad (53)$$

$\sigma_0(0) = 0$  and  $\Gamma_N = \emptyset$ . Then, if  $\text{supp}u_0$  is a non empty compact subset of  $\Omega$  the same happens with  $\text{supp}u(\cdot, t)$  for any  $t \in [0, t^*)$ , for some  $t^* \in (0, T]$ . Moreover, if  $t^* < T$  then  $u(x, t) > 0$  for any  $t \in (t^*, T]$ .

**Proof.** Consider  $w$  as the solution of the scalar homogeneous problem

$$\begin{cases} w_t - \varphi(w)_{xx} = 0 & \text{in } \Omega \times (0, T), \\ \varphi(w) = \varphi(u_D) \geq 0 & \text{on } \Gamma_D \times (0, T), \\ w(x, 0) = u_0(x) \geq 0 & \text{on } \Omega. \end{cases} \quad (54)$$

(remember that now  $\Gamma_D = \partial\Omega$ ). Thanks to the assumption (53) we know that there exists  $t^* \in (0, T]$  such that  $\text{supp}w(\cdot, t)$  is a compact subset of the open set  $\Omega$  for

any  $t \in [0, t^*)$  and that if  $t^* < T$  then  $w(x, t) > 0$  for any  $t \in (t^*, T]$ . It is easy to see that, necessarily,  $w(\cdot, t)$  must coincide with  $u(\cdot, t)$  for any  $t \in [0, t^*)$ . Indeed, as  $\sigma(u(x, t))v_x(x, t)$  must be a constant (in  $x$ )  $J(t)$  we get that, necessarily  $J(t) = 0$  if  $\sigma(u(x_0, t)) = 0$  for some  $x_0 \in \Omega$ . Then, as  $\sigma(w(x_0, t)) = 0$ , for some  $x_0 \in \Omega$  if  $t \in [0, t^*)$ , we can take  $v = z$  as the unique function solution of

$$\begin{cases} z_x = 0 & \text{in } \Omega \times (0, t^*), \\ z = v_D, & \text{on } \Gamma_D \times (0, t^*), \\ \frac{\partial z}{\partial n} = 0, & \text{on } \Gamma_N \times (0, t^*), \end{cases}$$

and we get that  $(w, z)$  satisfies problem (1) on  $[0, t^*)$  (notice that  $\sigma(w(x, t))(z_x(x, t))^2 = J(t)z_x(x, t) = 0$  on  $\Omega \times (0, t^*)$ ). By the uniqueness of solutions for problem (1) we conclude that  $(u, v) = (w, z)$  on  $\Omega \times [0, t^*)$ . Moreover, as  $\sigma(u(x, t))(v_x(x, t))^2 \geq 0$  on  $\Omega \times (0, T)$  we conclude (by the maximum principle for problem (54)) that  $u(x, t) \geq w(x, t) \geq 0$  on  $\Omega \times [0, T]$  and then  $u(x, t) > 0$  on  $\Omega \times (t^*, T]$ . ■

**Remark 4.** Notice that assumption (53) holds under the Wiedemann-Franz law  $k(u) = k_0 u \sigma(u)$  (remember (10)).

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