



On the confinement of a
viscous fluid by means of a
feedback external field

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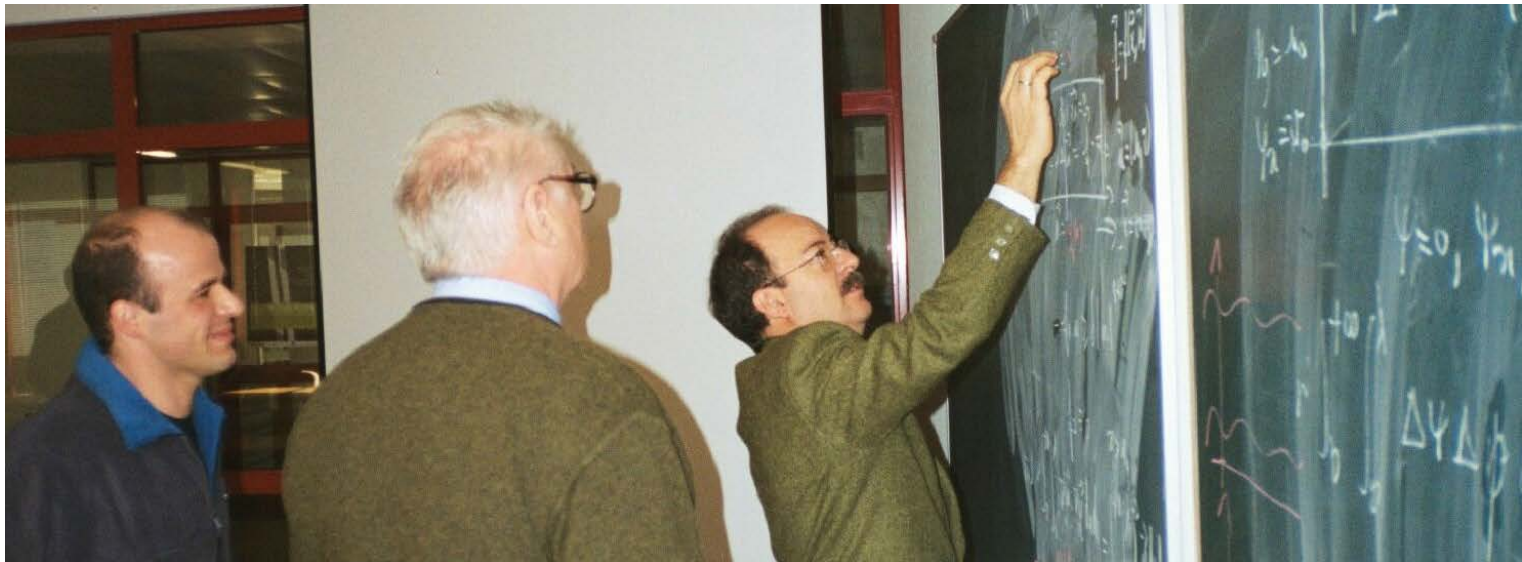


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planar stationary flow of an incompressible viscous fluid

$$\Omega = (0, \infty) \times (0, L), \quad \mathbf{u}(\mathbf{x}) = (u(\mathbf{x}), v(\mathbf{x})) \quad \mathbf{x} = (x, y) \in \Omega,$$

$$-\nu \Delta \mathbf{u} + (\mathbf{u} \cdot \nabla) \mathbf{u} = \mathbf{f} - \nabla p \quad \text{in } \Omega,$$

$$\operatorname{div} \mathbf{u} = 0 \quad \text{in } \Omega,$$

$$\mathbf{u}(0, y) = \mathbf{u}_*(y), \quad y \in (0, L)$$

$$\mathbf{u}(x, 0) = \mathbf{u}(x, L) = \mathbf{0}, \quad x \in (0, \infty).$$

$$|\mathbf{u}(x, y)| \rightarrow 0, \quad \text{as } x \rightarrow \infty \text{ and } y \in (0, L).$$

$\mathcal{P}(\Omega, \mathbf{u}_*, \mathbf{f})$

Problem: can we find an external localized

forces field \mathbf{f} stopping the fluid at a finite distance, i.e., such that

$$\mathbf{u}(x, y) = \mathbf{0} \quad \text{for } x > x_u \text{ and } y \in (0, L), \quad ?$$

$$\mathbf{f}(\mathbf{x}, \mathbf{u}) = \mathbf{0} \quad \text{for } x \geq x_f \text{ and } y \in (0, L),$$

compatibility conditions

$$\mathbf{u}_*(0) = \mathbf{u}_*(L) = \mathbf{0}.$$

$$\int_0^L u_*(s) ds = 0. \quad \mathbf{u}_*(y) = (u_*(y), v_*(y))$$

Remarks:

1. Resemblance with the question of the (magnetic) confinement of a plasma (ideal gas), typical of the magnetohydrodynamics (MHD).

2. For the classical Navier-Stokes problem ($\mathbf{f} \equiv \mathbf{f}(\mathbf{x})$) the decay is exponential and so the *localization effect* fails (Horgan and Wheeler, Amick, Ladyzhenskaya and Solonnikov, Ames and Payne, Iosif'yan: see book by Galdi (1994)). More recently: *Unique Continuation results*: L. Brandolese (CRAS, 2001), C. Fabre and G. Lebau (Comm. in PDE, 2002), and Y. Meyer (book in preparation).

Main result: Positive answer for $\mathbf{f}(\mathbf{x}) = \mathbf{f}(\mathbf{x}, \mathbf{u}(\mathbf{x}))$ (external body forces given in a feedback dissipative form)

$$\mathbf{f} : \Omega \times \mathbb{R}^2 \rightarrow \mathbb{R}^2, \quad \mathbf{f}(\mathbf{x}, \mathbf{u}) = (f_1(\mathbf{x}, \mathbf{u}), f_2(\mathbf{x}, \mathbf{u})),$$

and such that,

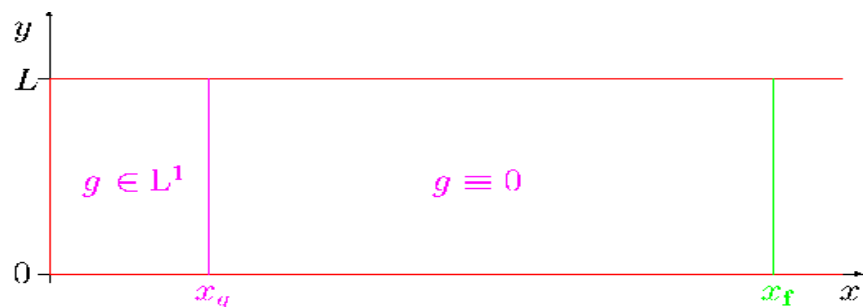
$$-\mathbf{f}(\mathbf{x}, \mathbf{u}) \cdot \mathbf{u} \geq \delta \chi_{\mathbf{f}}(\mathbf{x}) |u|^{1+\sigma} - g(\mathbf{x}) \quad (1)$$

$$\chi_{\mathbf{f}}(\mathbf{x}) := \chi_{(0, x_{\mathbf{f}})}(\mathbf{x}) = \begin{cases} 1 & \text{if } x \in (0, x_{\mathbf{f}}), \\ 0 & \text{if } x \notin (0, x_{\mathbf{f}}). \end{cases}$$

for some $0 < \sigma < 1$ and $\delta > 0$, $\forall \mathbf{u} \in \mathbb{R}^2$, $\mathbf{u} = (u, v)$, and a.e. $\mathbf{x} \in \Omega$, with

$$g \in L^1(\Omega^{x_g}), \quad g \geq 0, \quad g(\mathbf{x}) = 0 \quad \text{a.e. in } \Omega_{x_g} \quad (2)$$

for some x_g , with $0 \leq x_g \leq x_{\mathbf{f}} \leq \infty$ and $x_{\mathbf{f}}$ large enough, where $\Omega^{x_g} = (0, x_g) \times (0, L)$ and $\Omega_{x_g} = (x_g, \infty) \times (0, L)$.



Example: $\mathbf{f}(\mathbf{x}, \mathbf{u}) = -\delta \chi_{\mathbf{f}}(\mathbf{x}) (|u|^{\sigma-1} u, 0)$

Main idea to prove the localization property: introducing the current function ψ ,

$$u = \psi_y \quad \text{and} \quad v = -\psi_x \quad \text{in } \Omega, \quad (3)$$

we reduce $\mathcal{P}(\Omega, \mathbf{u}_*, \mathbf{f})$ to the scalar quasilinear fourth order problem (\mathcal{P}_ψ)

$\mathbf{f} = (f_1, f_2)$
 $\psi = 0, \quad \frac{\partial \psi}{\partial n} = 0$
 $\psi \rightarrow \infty$
 $|\nabla \psi| \rightarrow \infty$
 $\nu \Delta^2 \psi + \frac{\partial f_1}{\partial y}(\mathbf{x}, \psi_y, -\psi_x) - \frac{\partial f_2}{\partial x}(\mathbf{x}, \psi_y, -\psi_x)$
 $= \psi_y \Delta \psi_x - \psi_x \Delta \psi_y$
 $\psi = \int_0^y u_*(s) ds$
 $\frac{\partial \psi}{\partial n} = v_*(y)$
 $\psi = 0, \quad \frac{\partial \psi}{\partial n} = 0$

- Our starting point was a previous unidirectional result for second order anisotropic equations proved in S. N. Antontsev, J. I. Díaz and S. I. Shmarev, Birkhäuser, 2002, (Section 1.4.2).

- For the adaptation to four order equations we apply the technique introduced by F. Bernis (1984,...).

We assume

$$\mathbf{u}_* \in \mathbf{H}^{\frac{1}{2}}(0, L).$$

We introduce the functional spaces

$$\tilde{\mathbf{H}}(\Omega) = \{\mathbf{u} \in \mathbf{H}^1(\Omega) : \operatorname{div} \mathbf{u} = 0, \mathbf{u}(0, \cdot) = \mathbf{u}_*(\cdot), \mathbf{u}(x, L_1(x)) = \mathbf{u}(x, L_2(x)) = \mathbf{0}, x \in (0, \infty), \lim_{x \rightarrow +\infty} |\mathbf{u}| = 0\},$$

$$\tilde{\mathbf{H}}_0(\Omega) = \{\mathbf{u} \in \mathbf{H}^1(\Omega) : \operatorname{div} \mathbf{u} = 0, \mathbf{u}(0, \cdot) = \mathbf{0}, \mathbf{u}(x, L_1(x)) = \mathbf{u}(x, L_2(x)) = \mathbf{0}, x \in (0, \infty), \lim_{x \rightarrow \infty} |\mathbf{u}| = 0\}.$$

Assumptions for the existence of solutions:

$$\mathbf{f}(\mathbf{x}, \mathbf{u}) = -\delta \chi_{\mathbf{f}}(\mathbf{x}) (|\mathbf{u}|^{\sigma-1} \mathbf{u}(\mathbf{x}), 0) - \mathbf{h}(\mathbf{x}, \mathbf{u})$$

$$\mathbf{h}(\mathbf{x}, \mathbf{u}) \cdot \mathbf{u} \geq -g(\mathbf{x}) \text{ for every } \mathbf{u} \in \mathbb{R}^2 \text{ and a.e. } \mathbf{x} \in \Omega,$$

$$H_M(\mathbf{x}) = \sup_{|\mathbf{u}| \leq M} |\mathbf{h}(\mathbf{x}, \mathbf{u})| \quad H_M \in L^1(\Omega^{x_f}) \cap L^2(\Omega^{x_f}) \text{ for all } M > 0,$$

no upper restriction on the growth of $|\mathbf{f}(\mathbf{x}, \mathbf{u})|$ with respect to \mathbf{u} is imposed
this type of non-linear terms are called *strongly non linear*.

Definition. A function \mathbf{u} is a weak solution of problem $\mathcal{P}(\Omega, \mathbf{u}_*, \mathbf{f})$ if:

(i) $\mathbf{u} \in \tilde{\mathbf{H}}(\Omega)$, $\mathbf{f}(\mathbf{x}, \mathbf{u}) \in \mathbf{L}_{loc}^1(\Omega)$ and (ii) for every $\varphi \in \tilde{\mathbf{H}}_0(\Omega) \cap \mathbf{L}^\infty(\Omega)$ with compact support

$$\nu \int_{\Omega} \nabla \mathbf{u} : \nabla \varphi \, d\mathbf{x} + \int_{\Omega} (\mathbf{u} \cdot \nabla) \mathbf{u} \cdot \varphi \, d\mathbf{x} = \int_{\Omega} \mathbf{f} \cdot \varphi \, d\mathbf{x}.$$

Theorem 1. There exists, at least, one weak solution of problem $\mathcal{P}(\Omega, \mathbf{u}_*, \mathbf{f})$.

Moreover, $\mathbf{h}(\mathbf{x}, \mathbf{u}) \cdot \mathbf{u} \in L^1(\Omega)$, $\mathbf{h}(\mathbf{x}, \mathbf{u}) \in \mathbf{L}^1(\Omega)$ and \mathbf{u} satisfies to the energy estimate

$$\int_{\Omega} (|\nabla \mathbf{u}|^2 + \chi_{\mathbf{f}} |u|^{1+\sigma} + |\mathbf{h}(\mathbf{x}, \mathbf{u}) \cdot \mathbf{u}|) \, d\mathbf{x} \leq C \left(\|\mathbf{u}_*\|_{\mathbf{H}^{\frac{1}{2}}(L_1(0), L_2(0))}^4 + \|g\|_{L^1(\Omega^{x_g})} + 1 \right) \quad (4)$$

where $C = C(L, \delta, \nu, \sigma)$ is a positive constant. If, in addition, we assume

$$(\mathbf{f}(\mathbf{x}, \mathbf{u}_1) - \mathbf{f}(\mathbf{x}, \mathbf{u}_2)) \cdot (\mathbf{u}_1 - \mathbf{u}_2) \leq 0,$$

and ν large enough then $\mathcal{P}(\Omega, \mathbf{u}_*, \mathbf{f})$ has only one solution.

